# Probing Primordial Non-Gaussianity with Large-Scale Structure 

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## Primordial Non-Gaussianity from Inflation

Gaussianity is a consequence of:
i) inflaton a single scalar field
ii) slowly rolling
iii) in vacuum state
iv) with canonical kinetic terms
if we relax i) we have for the Bardeen potential,

$$
\Phi=\phi+f_{\mathrm{NL}} \phi^{2}
$$

which implies for it a bispectrum,

$$
B=2 f_{\mathrm{NL}} P_{1} P_{2}+\text { cyc. } \quad-10<f_{\mathrm{NL}}^{\text {local }}<74
$$

- For biased tracers (galaxies, halos), this model leads to a scale-dependent bias at large scales (Dalal et al 2008),

$$
b_{1}(k)=b_{10}+\Delta b_{1}\left(k, f_{\mathrm{NL}}\right)
$$

with $b \sim 1 / k^{\wedge} 2$ at low-k.Thus the power spectrum of galaxies is sensitive to fnl!!

## Generic Predictions in Peak-Background Split

We are interested in establishing as rigorously as possible the validity of the local PNG bias formula

$$
\Delta b_{1}\left(k, f_{\mathrm{NL}}\right)=\frac{2 f_{\mathrm{NL}}}{M(k)}\left(b_{10}-1\right) \delta_{c}
$$

and generalizing it to arbitrary (non-local) PNG. Some issues in derivations,

- proper treatment of filter and transfer function effects
- dependence on primordial bispectrum (cannot be just a number)
- peaks in phi vs peaks in delta approximations

$$
\nabla \phi^{2}=2 \phi \nabla^{2} \phi+2 \nabla \phi \cdot \nabla \phi \approx 2 \phi \nabla^{2} \phi ?
$$

simulations suggest a somewhat smaller amplitude (depending on halo def)
we know that "PBS" (in std implementation) is not that accurate even in Gaussian case:

link 0.2

link 0.2

Manera, Sheth, R.S. (2010)

Figure 11. Comparison of large-scale bias estimates for the same halo mass bins as in previous figures when $l_{\text {link }}=0.2$. Thick bars show the measured $P_{\mathrm{hm}} / P_{\mathrm{mm}}$ (left-hand panel) and $\sqrt{\xi_{\mathrm{hh}} / \xi_{\mathrm{mm}}}$ (right-hand panel), and symbols with error bars show the linear bias parameter $b_{1}$ predicted from the peakbackground split.

Insight into mass function and bias is provided by the excursion-set formalism of halo formation (Bond et al I991).


A full calculation of the PBS change in bias due to arbitrary PNG bispectrum gives,

$$
\begin{gathered}
\Delta b(k)=\frac{\partial_{\sigma^{2}}\left[I_{B}(k) \mathcal{F}_{0}\right]}{M(k) \mathcal{F}_{0}} \\
I_{B}(k, R) \equiv \frac{1}{P_{\phi}(k)} \int B_{\delta_{R} \delta_{R} \phi}(\boldsymbol{q}, \boldsymbol{k}-\boldsymbol{q},-\boldsymbol{k}) d^{3} q
\end{gathered}
$$

Note that, unlike the GW86 formula, what matters is the *cross* bispectrum. For local PNG, expanding in powers of $k$ small (with higher-order corrections coming from filter, transfer function, grad-phi terms, etc

$$
I_{B}(k=0, R) \approx 4 f_{\mathrm{NL}} \sigma_{R}^{2}(m)+\mathcal{O}\left(k^{2}\right)
$$

which gives
non-markovian

$$
\Delta b(k)=\frac{4 f_{\mathrm{NL}}}{M(k)} \partial_{\ln \sigma^{2}} \ln \left(\sigma^{2} \mathcal{F}_{0}\right) \stackrel{\downarrow}{<} \frac{2 f_{\mathrm{NL}}}{M(k)} \delta_{c} \frac{\left(\partial \mathcal{F} / \partial \delta_{\ell}\right)_{0}}{\mathcal{F}_{0}}=\frac{2 f_{\mathrm{NL}}}{M(k)} \delta_{c}\left(b_{1}-1\right)
$$

the precise relationship has to be obtained from the first-crossing prob FO.




## Large Suite of Dark Matter Simulations (LasDamas)



Oriana
with
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## LasDamas Simulations

| Name | Sample | Lbox | Npar | mpar | Nrealiz |
| :---: | :---: | :---: | :---: | :---: | :---: |
| Oriana (G) | $\begin{gathered} \text { LRG } \\ + \text { Main - } 22 \end{gathered}$ | 2400 | 1280^3 | 4.57E+1I | 42 |
| Oriana <br> fnl_local=+100 | $\begin{gathered} \text { LRG } \\ + \text { Main - } 22 \end{gathered}$ | 2400 | I280^3 | 4.57E+II | 12 |
| $\begin{gathered} \text { Oriana } \\ \text { fnl_equi }=-400 \end{gathered}$ | $\begin{gathered} \text { LRG } \\ + \text { Main - } 22 \end{gathered}$ | 2400 | I280^3 | 4.57E+II | 12 |
| Oriana <br> fnl_orto=-400 | $\begin{gathered} \text { LRG } \\ + \text { Main - } 22 \end{gathered}$ | 2400 | I280^3 | 4.57E+II | 12 |
| Carmen | Main -2I | 1000 | 1120^3 | $4.98 \mathrm{E}+10$ | 42 |
| Esmeralda | Main -20 | 640 | 1250^3 | $9.31 \mathrm{E}+09$ | 50 |
| Consuelo | Main -19-18 | 420 | 1400^3 | 1.87E+09 | 50 |

Nmocks=4 x Nrealiz, 2LPT ICs, Gaussian Mocks available at http:///ss.phy.vanderbilt.edu/lasdamas/

## Large-Scale Bias in non-local PNG: Simulations

- In single-field inflationary models, we are instead interested in models that correspond to non-local PNG (due to non-canonical kinetic terms). For example, the equilateral model has a Bardeen potential bispectrum,

$$
\begin{aligned}
\left(6 f_{\mathrm{NL}}\right)^{-1} B_{\text {equil }}=-P_{1} P_{2}-2\left(P_{1} P_{2} P_{3}\right)^{2 / 3}+ & P_{1}^{1 / 3} P_{2}^{2 / 3} P_{3} \\
& -214<f_{\mathrm{NL}}^{\text {equil }}<266
\end{aligned}
$$

(permutations are understood), whereas the orthogonal model template reads

$$
\begin{array}{r}
\left(6 f_{\mathrm{NL}}\right)^{-1} B_{\text {ortho }}=-3 P_{1} P_{2}-8\left(P_{1} P_{2} P_{3}\right)^{2 / 3}+3 P_{1}^{1 / 3} P_{2}^{2 / 3} P_{3} \\
-410<f_{\mathrm{NL}}^{\text {ortho }}<6
\end{array}
$$

We are interested in generating such bispectra from quadratic (non-local) models, i.e.

$$
\Phi=\phi+f_{\mathrm{NL}} K[\phi, \phi]
$$

where K is the appropriate non-local quadratic kernel that generates the desired bispectrum. For simplicity, here we assume scale-invariance.

- Introduce some handy non-local operators

$$
\begin{gathered}
\partial \phi \equiv \sqrt{-\nabla^{2}} \phi(\mathbf{x}) \equiv \int \mathrm{e}^{-i \mathbf{k} \cdot \mathbf{x}} k \phi(\mathbf{k}) d^{3} k \\
\nabla^{-2} A(\mathbf{x}) \equiv-\int \mathrm{e}^{-i \mathbf{k} \cdot \mathbf{x}}\left(\frac{1}{k^{2}}\right) A(\mathbf{k}) d^{3} k \\
\partial^{-1} A \equiv \sqrt{-\nabla^{-2}} A \equiv \int \mathrm{e}^{-i \mathbf{k} \cdot \mathbf{x}}\left(\frac{1}{k}\right) A(\mathbf{k}) d^{3} k
\end{gathered}
$$

Then the EQ and ORT bispectra templates can be generated by,
$K[\phi, \phi]=a \phi^{2}+b \partial^{-1}(\phi \partial \phi)+c \nabla^{-2}\left(\phi \nabla^{2} \phi\right)+d \nabla^{-2}(\partial \phi)^{2}+e \nabla^{-2} \partial^{-1}\left(\phi \nabla^{2} \partial \phi\right)+f \nabla^{-2} \partial^{-1}\left(\nabla^{2} \phi \partial \phi\right)$
regularity constraints (one-loop corrections to the power spectrum must preserve scale-invariance in the IR) restrict the free parameters that leave the bispectrum invariant. Note these kernels have correct exchange symmetry.




CrossBias from Oriana Simulations $z=0,0.34,0.97$

$P_{\mathrm{h}}^{\mathrm{ZA}} / P_{\mathrm{h}}$





Halo $\mathrm{S} / \mathrm{N}$ for Non-Gaussian Models, $\mathrm{z}=1.0, \mathrm{M}>10^{14} \mathrm{Mo}$
dashed $=$ from bispectrum, solid=from power


Adding Bispectrum information helps a lot... Signal to Noise $\mathrm{f}_{\mathrm{NL}}=100$ (local)
LRG mocks including redshift distortions, Mag <21.2, $\mathrm{z}=0.342$


## Things to worry about when using Bisp for PNG...

Shapes induced by:

- gravitational instability (2-loop RPT for the mass)
- Redshift distortions beyond PT (extension of RS 04)
- Galaxy bias beyond local approximation (even for Gaussian ICs)
- characterize Bispectrum Eigenmodes + non-Gaussian likelihood (RS 00; Gaztanaga \& RS 05)


## Bispectrum



